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## LETTER TO THE EDITOR

## Universal $\boldsymbol{R}$-matrix for a two-parametric quantization of $\mathbf{g}(\mathbf{2})$

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#### Abstract

For the two-parametric quantized group $\mathrm{gl}_{4, s}(2)$, we calculate the explicit form of the universal $R$-matrix.


The deformations of the algebra $\mathrm{gl}_{q, s}(2)$ with two deformation parameters have been studied in many papers (Kulish 1990, Schirrmacher et al 1991, Burdík and Hlavatý 1991, Schirrmacher 1991). The starting point for a construction of this two-parametric deformation is a two-parametric solution of the Yang-Baxter equation (ybe)

$$
\begin{equation*}
R_{12} R_{13} R_{23}=R_{23} R_{13} R_{12} \tag{1}
\end{equation*}
$$

found by Hlavatý (1987), Demidov et al (1990) and Schirrmacher et al (1991):

$$
R_{q, s}=\left(\begin{array}{cccc}
q & 0 & 0 & 0 \\
0 & s & 0 & 0 \\
0 & \Omega & 1 / s & 0 \\
0 & 0 & 0 & q
\end{array}\right) q^{-1 / 2}
$$

where

$$
\begin{equation*}
\Omega=q-q^{-1} . \tag{2}
\end{equation*}
$$

It was found by Reshetikhin, Takhtajan and Faddeev (1989) that the defining relations for the quantum enveloping algebra can be obtained from the relations

$$
\begin{align*}
& R^{+} L_{1}(\varepsilon) L_{2}(\varepsilon)=L_{2}(\varepsilon) L_{1}(\varepsilon) R^{+} \\
& R^{+} L_{1}(+) L_{2}(-)=L_{2}(-) L_{1}(+) R^{+} \tag{3}
\end{align*}
$$

where

$$
\begin{align*}
& R^{+}:=P R P  \tag{4}\\
& P:=\left(\begin{array}{llll}
1 & 0 & 0 & 0 \\
0 & 0 & 1 & 0 \\
0 & 1 & 0 & 0 \\
0 & 0 & 0 & 1
\end{array}\right) \\
& \varepsilon= \pm \quad L_{1}(\varepsilon)=L(\varepsilon) \otimes 1 \quad L_{2}(\varepsilon)=1 \otimes L(\varepsilon) \tag{5}
\end{align*}
$$

The explicit calculation for the $R$-matrix (1) was given by Kulish (1990). In this case

$$
L(+)=\left(\begin{array}{cc}
k_{1} & Y  \tag{6}\\
0 & k_{2}
\end{array}\right) \quad L(-)=\left(\begin{array}{cc}
l_{1} & 0 \\
X & l_{2}
\end{array}\right) \quad \text { and } R=R_{q, s}
$$

The algebra $\mathrm{gl}_{q, s} \equiv \mathrm{U}_{q, s}(\mathrm{gl}(2))$ is obtained as the free algebra generated by ( $k_{1}, k_{2}, l_{1}, l_{2}, X, Y$ ) and factorized by the ideal corresponding to the relations

$$
\begin{array}{lcc}
k_{1} Y=(q s)^{-1} Y k_{1} & k_{2} Y=q s^{-1} Y k_{2} & {\left[k_{1}, k_{2}\right]=0} \\
l_{1} X=(q)^{-1} s X l_{1} & l_{2} X=q s X l_{2} & {\left[l_{1}, l_{2}\right]=0} \\
k_{1} X=q s X k_{1} & k_{2} X=q^{-1} s X k_{2} & {\left[k_{i}, l_{j}\right]=0}  \tag{7}\\
l_{1} Y=q s^{-1} Y l_{1} & l_{2} Y=(q s)^{-1} X l_{2} & \\
{[X, Y]_{s}=s X Y-1 / s Y X=\Omega\left(k_{2} l_{1}-l_{2} k_{1}\right)} &
\end{array}
$$

The familiar commutation relations of the $\mathrm{gl}_{q, s}(2)$

$$
\begin{array}{ll}
{\left[J_{0}, J_{ \pm}\right]= \pm J_{ \pm}} & {\left[J_{+}, J_{-}\right]=\frac{\sin \left(2 \eta J_{0}\right)}{\sinh \eta}}  \tag{8}\\
{\left[Z, J_{ \pm}\right]=0} & {\left[Z, J_{0}\right]=0}
\end{array}
$$

can be obtained from (7) by the substitutions

$$
\begin{array}{lll}
k_{1}=(q s)^{-J_{0}+Z} & k_{2}=\left(\frac{s}{q}\right)^{-J_{0}+Z} & Y=\Omega J_{+}(s)^{-J_{0}}  \tag{9}\\
l_{1}=\left(\frac{s}{q}\right)^{-j_{0}-Z} & l_{2}=(q s)^{-J_{0}-Z} & X=-\Omega J_{-}(s)^{-J_{0}}
\end{array}
$$

We can see that the above commutation relations (8) are not deformed with $s$ and this deformation parameter is only in the comultiplication for this algebra

$$
\begin{align*}
& \dot{\Delta}\left(J_{+}\right)=q^{-J_{0}}(q s)^{z} \otimes J_{+}+J_{+} \otimes q^{J_{0}}\left(\frac{s}{q}\right)^{z} \\
& \Delta\left(J_{-}\right)=q^{-J_{0}}(q s)^{-z} \otimes J_{-}+J_{-} \otimes q^{J_{0}}\left(\frac{s}{q}\right)^{-z}  \tag{10}\\
& \Delta\left(J_{0}\right)=J_{0} \otimes 1+1 \otimes J_{0}, \Delta(Z)=Z \otimes 1+1 \otimes Z
\end{align*}
$$

and in the antipod

$$
\begin{equation*}
S(Z)=-Z \quad S\left(J_{0}\right)=-J_{0} \quad S\left(J_{ \pm}\right)=-q^{ \pm 1} s^{\mp 2 Z} J_{ \pm} \tag{11}
\end{equation*}
$$

For a co-unit we have again

$$
\begin{equation*}
\varepsilon(Z)=\varepsilon\left(J_{0}\right)=\varepsilon\left(J_{ \pm}\right)=0 \tag{12}
\end{equation*}
$$

Theorem 1. The algebra $\mathrm{gl}_{q, s}(2)$ with $\Delta, \varepsilon$ and $S$ given by (10), (12) and (11) is a Hopf algebra.

Proof. By direct verification of the axioms of the Hopf algebra.
A Hopf algebra is a quasitriangular Hopf algebra (see Drinfeld 1986) if the comultiplications $\Delta, \Delta^{\prime}$ are related by conjugation

$$
\begin{equation*}
\Delta^{\prime}(a) \equiv \sigma \quad \Delta(a)=R \Delta(a) R^{-1} \quad R \in A \otimes A \quad \sigma(a \otimes b)=b \otimes a \tag{13}
\end{equation*}
$$

for any $a \in A$, and the following conditions are satisfied

$$
\begin{align*}
& (i d \otimes \Delta)(R)=R_{13} R_{12} \\
& (\Delta \otimes i d)(R)=R_{13} R_{23}  \tag{14}\\
& (S \otimes i d)(R)=R^{-1}
\end{align*}
$$

For the $\operatorname{sl}_{q}(2)$ the universal element

$$
\begin{equation*}
R=q^{\left(J_{0} \otimes J_{0}\right) / 2} \sum_{n=0}^{\infty} \frac{\left(1-q^{-2}\right)^{n}}{[n]_{q}!}\left(q^{J_{0} / 2} J_{+} \otimes q^{-J_{0} / 2} J_{-}\right)^{n} q^{n(n-1) / 2} \tag{15}
\end{equation*}
$$

was given by Kirillov and Reshetikhin (1988) and one can check it satisfies (13) and (14).
An explicit form of universal element is given in the following theorem.
Theorem 2. The algebra $\operatorname{gl}_{q, s}(2)$ is the quasitriangular Hopf algebra with the universal element $\bar{R}_{q, s} \in \mathrm{gl}_{q, s}(2) \otimes \mathrm{gl}_{q, s}(2)$ which is given by the formula

$$
\begin{gather*}
R_{q, s}=q^{2\left(J_{0} \otimes J_{0}+J_{0} \otimes z-z \otimes J_{0}\right)} \sum_{n=0}^{\infty} \frac{\left(1-q^{-2}\right)^{n}}{[n]_{q}!} q^{n(n-1) / 2}\left(q^{J_{0}}(q s)^{-z} J_{+}\right)^{n} \\
\otimes\left(q^{-J_{0}}\left(\frac{s}{q}\right)^{z} J_{-}\right)^{n} \tag{16}
\end{gather*}
$$

Proof. By direct verification of the relation in (13) and (14).
In this letter we have presented the universal element $R$ of the quantum group $\mathrm{gl}_{q . s}(2)$. We propose that this explicit realization could play a role in nonlinear integrable systems in physics.

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